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Variational calculus

Example of a problem in variational calculus:

Determine the shape of a soap bubble surface between two circular rings.

Solution: Choose cylindrical coordinates ρ, φ, z . The surface is described by $\rho = \rho(\varphi, z) = \rho(z)$. The rings sit at $z = z_1, \rho = a$ and $z = z_2, \rho = a$. The physical principle: area is minimized.

$$A[\rho] = \int_{z_1}^{z_2} dz \, 2\pi\rho(z) \sqrt{1 + (\rho'(z))^2}$$

with boundary conditions $\rho(z_1) = a, \rho(z_2) = a$. Variational calculus is a method to solve problems like: Given function F of three variables, find a function $\rho: [z_1, z_2] \rightarrow \mathbb{R}$ that minimizes

$$A[\rho] \equiv \int_{z_1}^{z_2} F(\rho(z), \rho'(z), z) dz$$

with given boundary conditions $\rho(z_1) = \rho_1, \rho(z_2) = \rho_2$.

Compare with ordinary differential calculus: in the similar problem a function $f(x_1, \dots, x_n)$ is minimized. The correspondence is

$$\begin{aligned} A &\leftrightarrow f \\ \rho(z) &\leftrightarrow x_i \\ z &\leftrightarrow i \\ [z_1, z_2] &\leftrightarrow \{1, \dots, n\} \end{aligned}$$

Reduce the problem to differential calculus as follows:

Suppose $\rho_0(z)$ gives the minimum. Consider another function $\rho(z) = \rho_0(z) + \Delta\rho(z)$ which satisfies the boundary conditions, but is otherwise arbitrary.

Consider the set of functions $\rho_\epsilon(z) = \rho_0 + \epsilon \Delta\rho$, where ϵ ranges over an interval around zero. $A(\epsilon) = A[\rho_\epsilon]$ has a minimum at $\epsilon = 0$.

$$\left. \frac{dA(\epsilon)}{d\epsilon} \right|_{\epsilon=0} = 0$$

$$A(\epsilon) - A(0) = \int_{z_1}^{z_2} dz \left(\frac{\partial F}{\partial \rho_0} \epsilon \Delta\rho + \frac{\partial F}{\partial \rho'} \epsilon \frac{d}{dz}(\Delta\rho) \right) + \mathcal{O}(\epsilon^2)$$

$$= \epsilon \left[\frac{\partial F}{\partial \rho'_0} \Delta\rho \right]_{z_1}^{z_2} + \epsilon \int_{z_1}^{z_2} dz \left(\frac{\partial F}{\partial \rho_0} - \frac{d}{dz} \frac{\partial F}{\partial \rho'} \right) \Delta\rho + \mathcal{O}(\epsilon^2) =$$

$\Delta\rho(z_1) = 0, \Delta\rho(z_2) = 0$:

$$= \epsilon \int_{z_1}^{z_2} dz \left(\frac{\partial F}{\partial \rho_0} - \frac{d}{dz} \frac{\partial F}{\partial \rho'} \right) \Delta\rho + \mathcal{O}(\epsilon^2)$$

$$0 = \left. \frac{\partial A(\epsilon)}{\partial \epsilon} \right|_{\epsilon=0} = \int_{z_1}^{z_2} dz \left(\frac{\partial F}{\partial \rho_0} - \frac{d}{dz} \frac{\partial F}{\partial \rho'} \right) \Delta\rho$$

$$\frac{\partial F}{\partial \rho_0} - \frac{d}{dz} \frac{\partial F}{\partial \rho'} = 0 \text{ for } z \in [z_1, z_2].$$

Same as Lagrange's equations in mechanics, with $F \rightarrow L, \rho \rightarrow q, z \rightarrow t$.

So, Lagrange's equations \Rightarrow **Hamilton's principle**:

A mechanical system moves between two times such that the **action** is stationary under small variations of the path such that the endpoints remain fixed:

$$\text{Action} = \int_{t_1}^{t_2} L(q, \dot{q}, t) dt$$

Note: Usually the stationary point is a saddle point (neither max nor min).

Note: One advantage of Hamilton's principle: It's a simply stated complete characterisation of possible motions of a system, manifestly independent of the choice of coordinate.

Note: In practice, in variational calculations, one is more brief: One writes

$$\delta A[q] = \delta \int_{t_1}^{t_2} L(q, \dot{q}) dt = \int \left(\frac{\partial L}{\partial q} \delta q + \frac{\partial L}{\partial \dot{q}} \delta \dot{q} \right) dt = \left[\frac{\partial L}{\partial \dot{q}} \delta q \right]_{t_1}^{t_2} + \int \left(\frac{\partial L}{\partial q} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} \right) \delta q dt$$

Note: System with holonomic constraints:

$$\frac{\partial L}{\partial q^\nu} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}^\nu} + \sum_{\alpha} \lambda_{\alpha}(t) \frac{\partial f_{\alpha}}{\partial q^{\nu}} = 0, \quad \text{constraints } f_{\alpha}(q, t) = 0$$

Described by: action

$$A[q, \lambda] = \int \left(L(q, \dot{q}) + \sum_{\alpha} \lambda_{\alpha} f_{\alpha}(q) \right) dt$$

Here $\lambda = \lambda(t)$ is a function of time.

Remark: In variational calculus Lagrange parameters which are not functions of time can also occur.

Example: Homogeneous chain. $\langle \dots \rangle$

Note: There do exist velocity dependent holonomic constraints, i.e. constraints $f_{\alpha}(q, \dot{q})$ which can be treated by forming L^{ext} .

I know no general method to decide from looking at a time dependent constraint whether it is holonomic or not.

Example illustrating the importance of boundary conditions.

Determine the shape of an elastic homogeneous rod fastened at its ends.

Solution:

Physics: potential energy is minimized. Describe shape by function $y(x), 0 \leq x \leq l$.

$$A[y] = V_{\text{gravitational}} + V_{\text{elastic}} = \int_0^l dx \left[\rho g y(x) + \frac{1}{2} \kappa \left(\frac{dy}{dx} \right)^2 \right]$$

(if the rod is approximately horizontal). Compare Hook's law and the energy of a spring. ρ is mass per unit length, κ is the stiffness of the rod.

$$\delta A = \int_0^l (\rho g \delta y + \kappa y'' \delta y')$$

$$\kappa y'' \delta y' = \kappa \frac{d}{dx} (y'' \delta y') - \kappa y''' \delta y'$$

$$-\kappa y''' \delta y' = -\frac{d}{dx} (\kappa y''' \delta y) + \kappa y^{\text{IV}} \delta y$$

$$\delta A = \kappa [y'' \delta y' - y''' \delta y]_0^l + \int_0^l dx (\rho g + \kappa y^{\text{IV}}) \delta y = 0$$

$\Rightarrow 1)$ Eqn: $pg + \kappa y^{\text{IV}} = 0, 0 \leq x \leq l$

2) Boundary conditions.

a) $\delta y(0) = 0, \delta y(l) = 0$, but $\delta y'$ arbitrary. The boundary terms in δA require that $y''(0) = 0, y''(l) = 0$. \Rightarrow in total 4 boundary conditions. Good, because the equation is of fourth order. We get four integration constants.

b) $\delta y(0) = 0, \delta y'(0) = 0$. No restriction on the other end. Boundary terms require $y''(l) = 0, y'''(l) = 0$. Again, in total 4 boundary conditions.